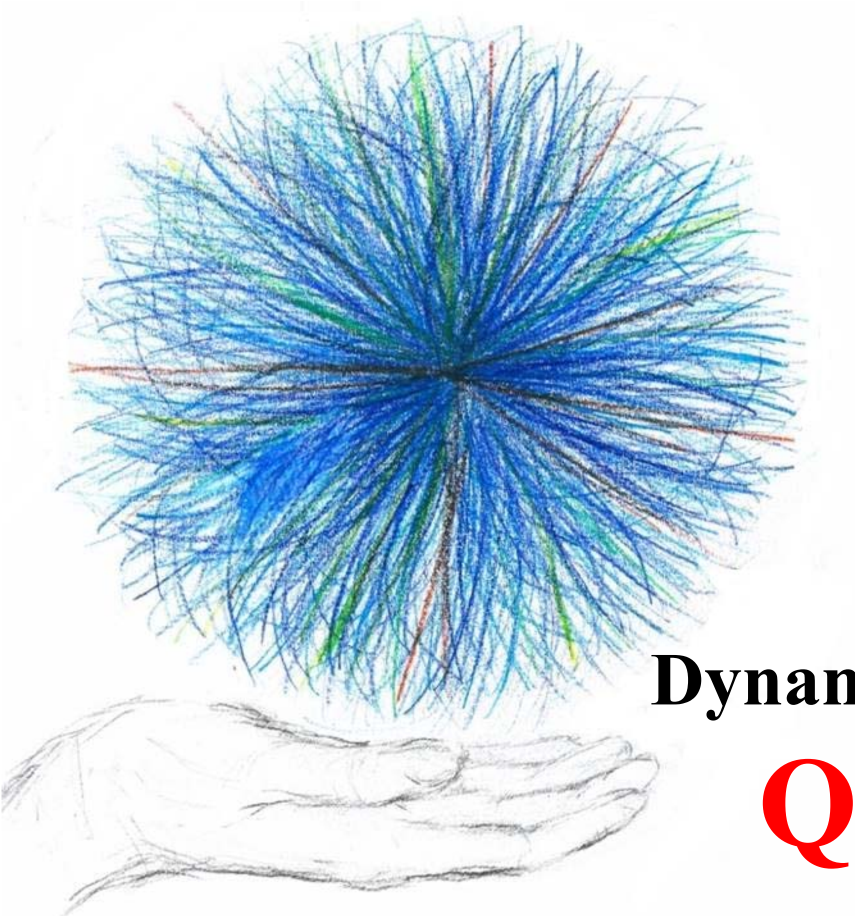


Dynamical energy loss as a tool for

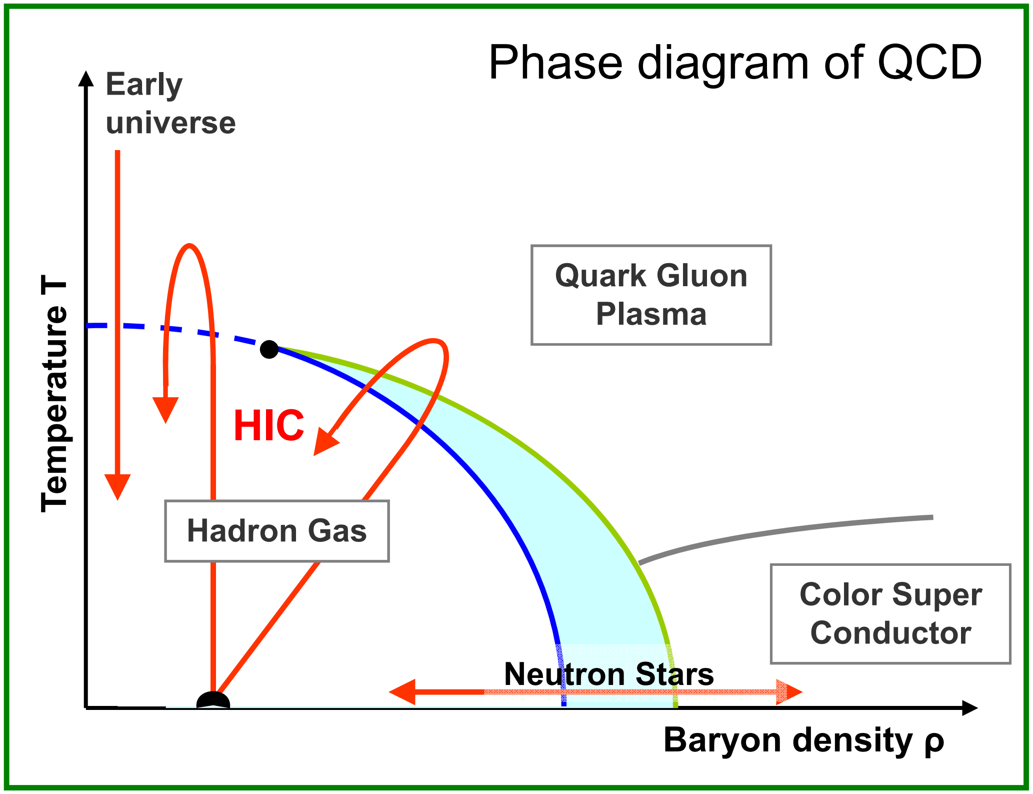
QGP Tomography

Magdalena Djordjevic, IPB

Brief overview of Quark Gluon Plasma

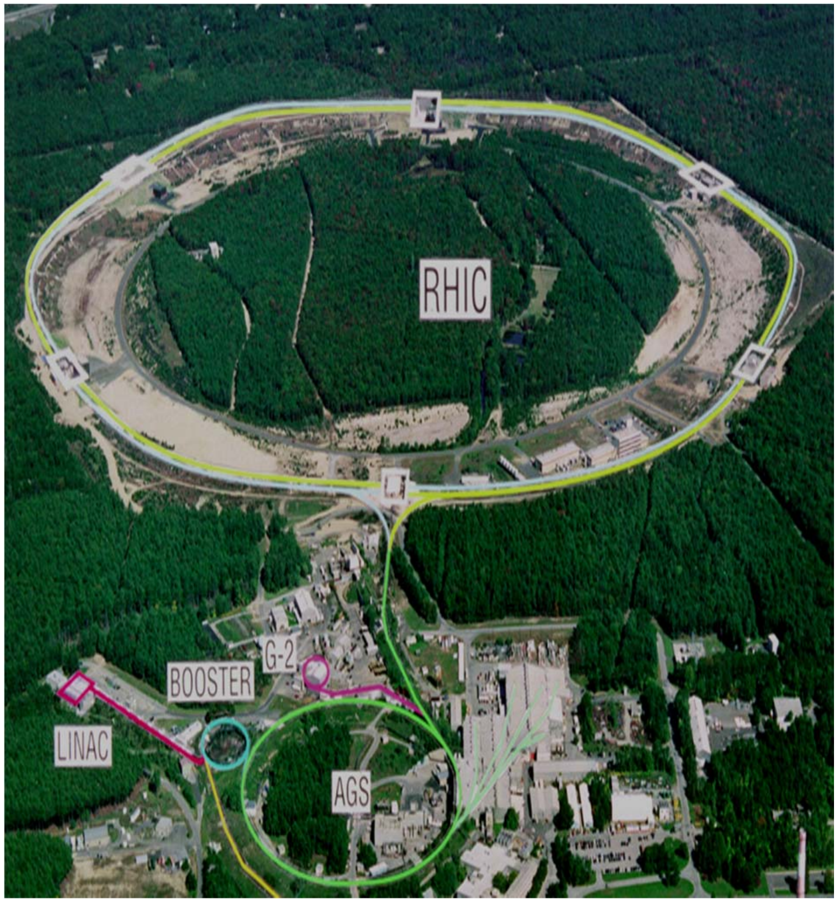
➢ QGP is a new form of matter, consisting of deconfined and interacting quarks, antiquarks and gluons.

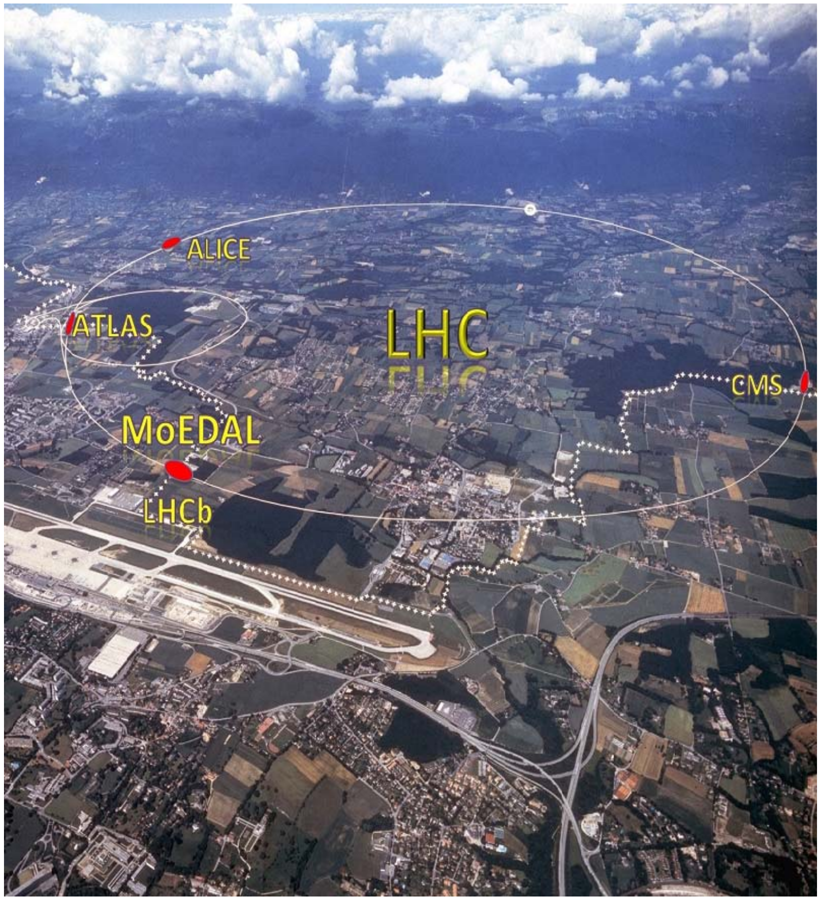
➢ QGP is predicted by QCD to exist at extremely high energy densities.



One of the most important goals of high energy heavy ion physics is to form, observe and understand QGP.

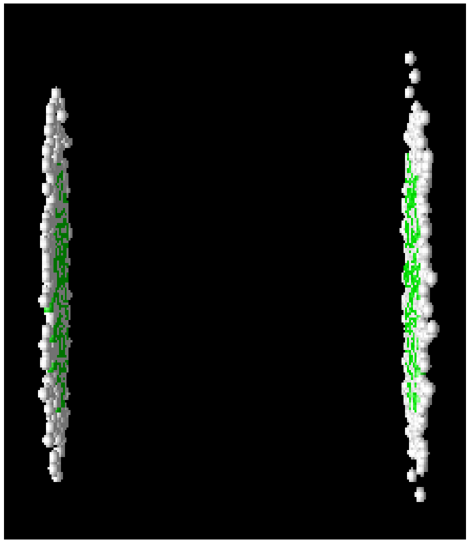
Ultra-Relativistic Heavy Ion Colliders (RHIC and LHC) have been made at BNL and CERN.



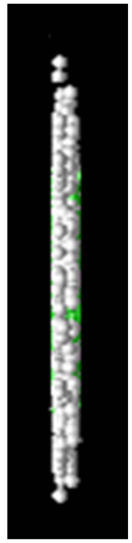


Scheme of relativistic heavy ion collisions

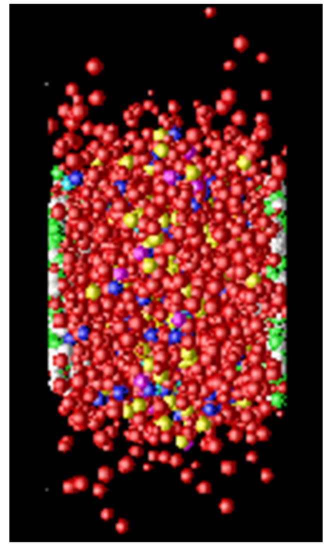
Simulation "VNI" (Geiger, Longacre, Srivastava)



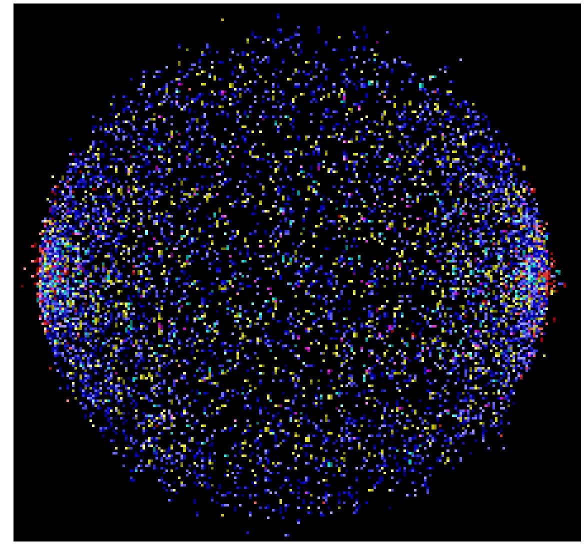
Heavy ion acceleration



Collision



Quark-gluon plasma



Hadron Gas

To study the properties of QCD matter created at URHIC we need good probes

High energy particles (E > 10 GeV) are widely recognized as the excellent probes of QGP.

Why are high energy particles good probes?

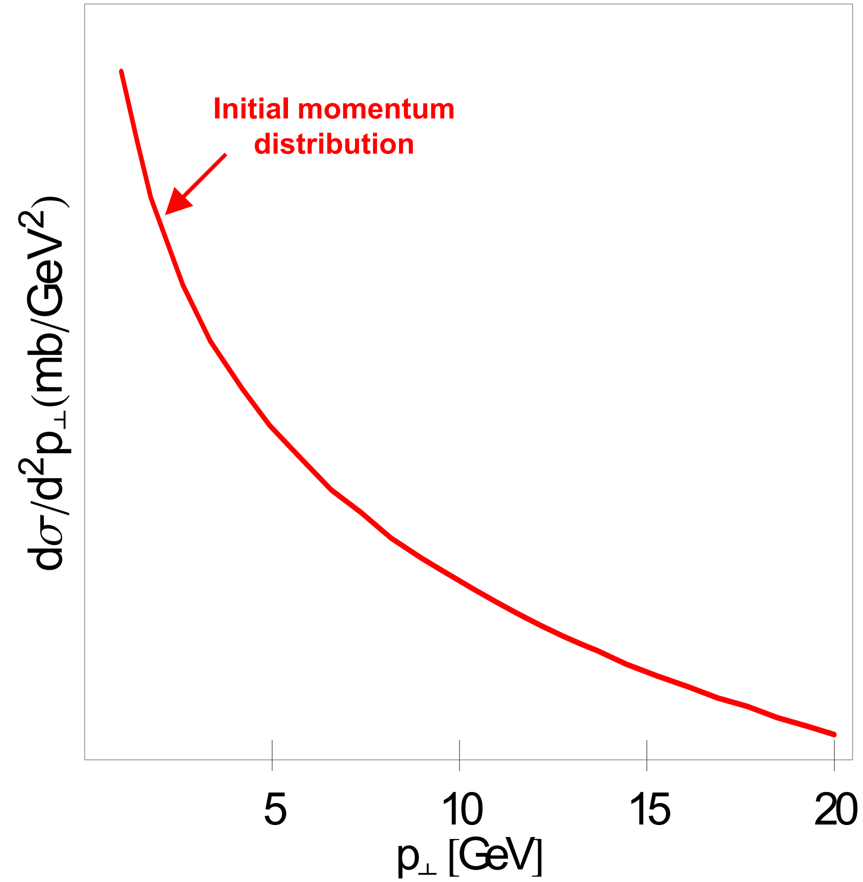
High energy particles:

- Are produced only during the early stage of QCD matter.
- Significantly interact with the QCD medium
- Perturbative calculations are possible

Jet suppression

Jet suppression is considered to be an excellent probe of QCD matter.

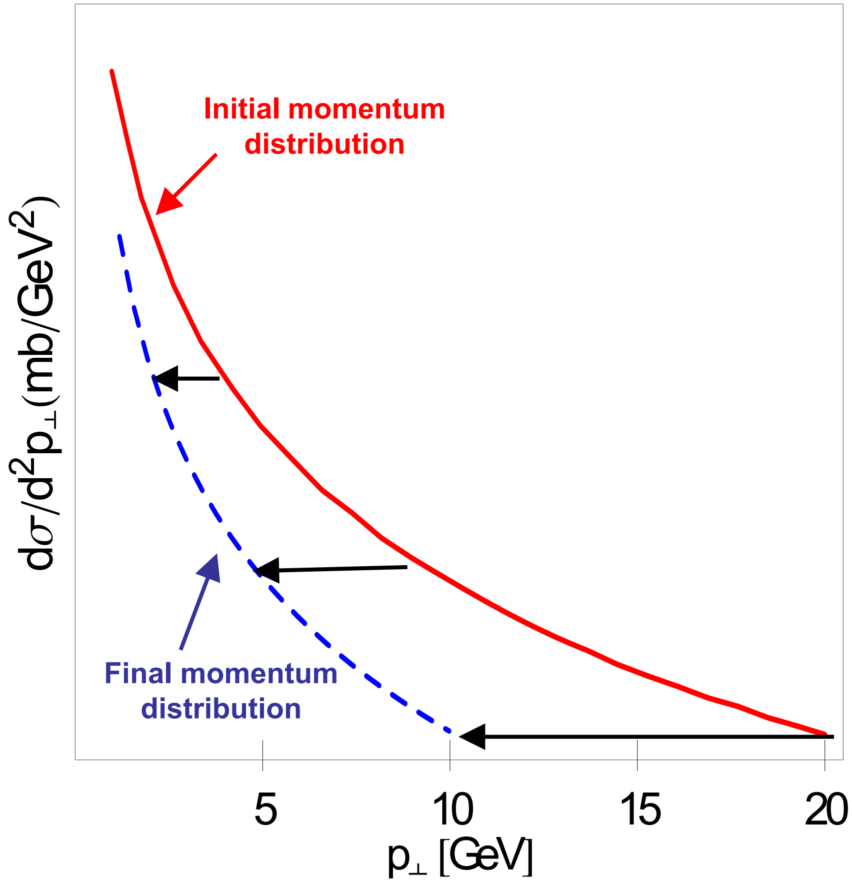
What is suppression?



Jet suppression

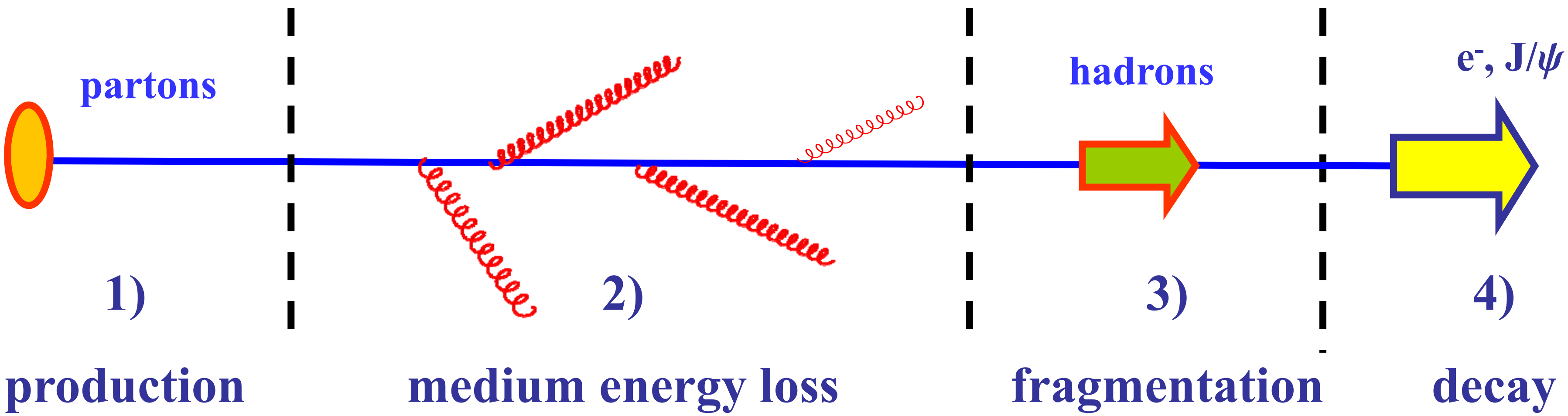
Jet suppression is considered to be an excellent probe of QCD matter.

What is suppression?



$$\text{Suppression} = \frac{\text{Final momentum distribution}}{\text{Initial momentum distribution}}$$

Suppression scheme

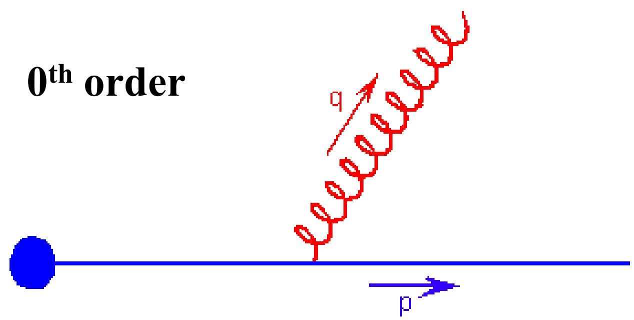


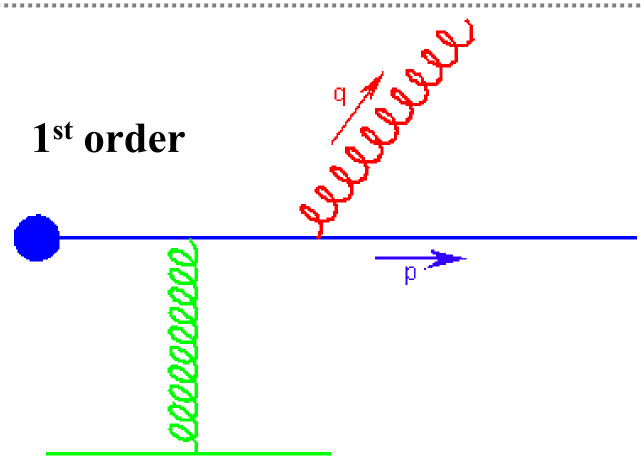
1) Initial momentum distributions for partons
2) Parton energy loss
3) Fragmentation functions of partons into hadrons
4) Decay of heavy mesons to single e$^-$ and J/ψ.

Energy loss in QGP

Radiative energy loss

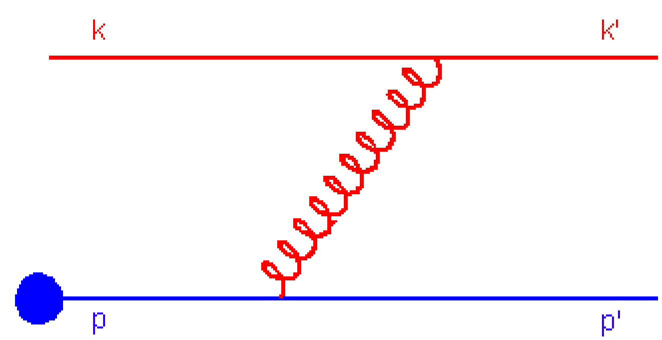
Radiative energy loss comes from the processes in which there are more outgoing than incoming particles:





Collisional energy loss

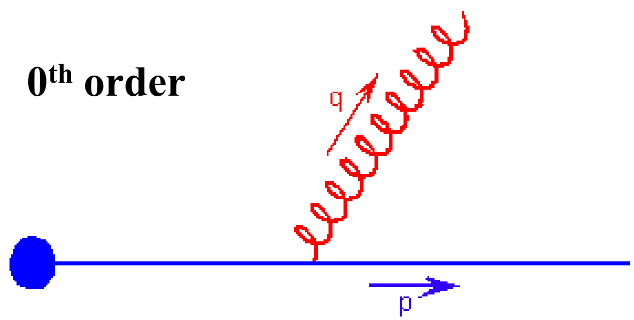
Collisional energy loss comes from the processes which have the same number of incoming and outgoing particles:



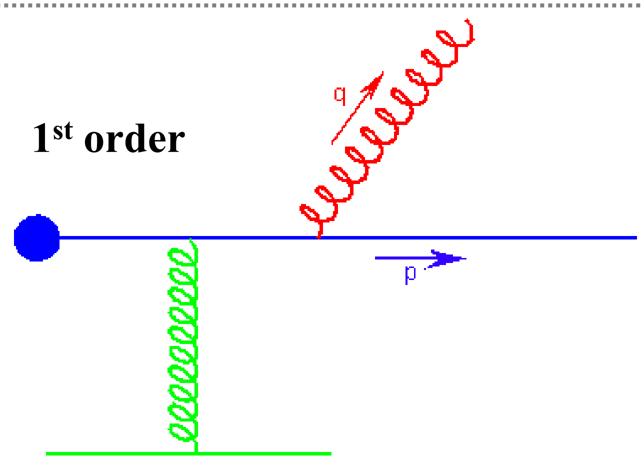
Radiative energy loss

Radiative energy loss comes from the processes in which there are more outgoing than incoming particles:

0^{th} order



1^{st} order



Collisional energy loss

Collisional energy loss comes from the processes which have the same number of incoming and outgoing particles:

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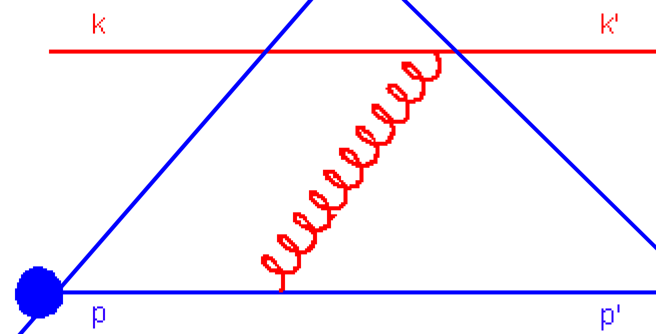
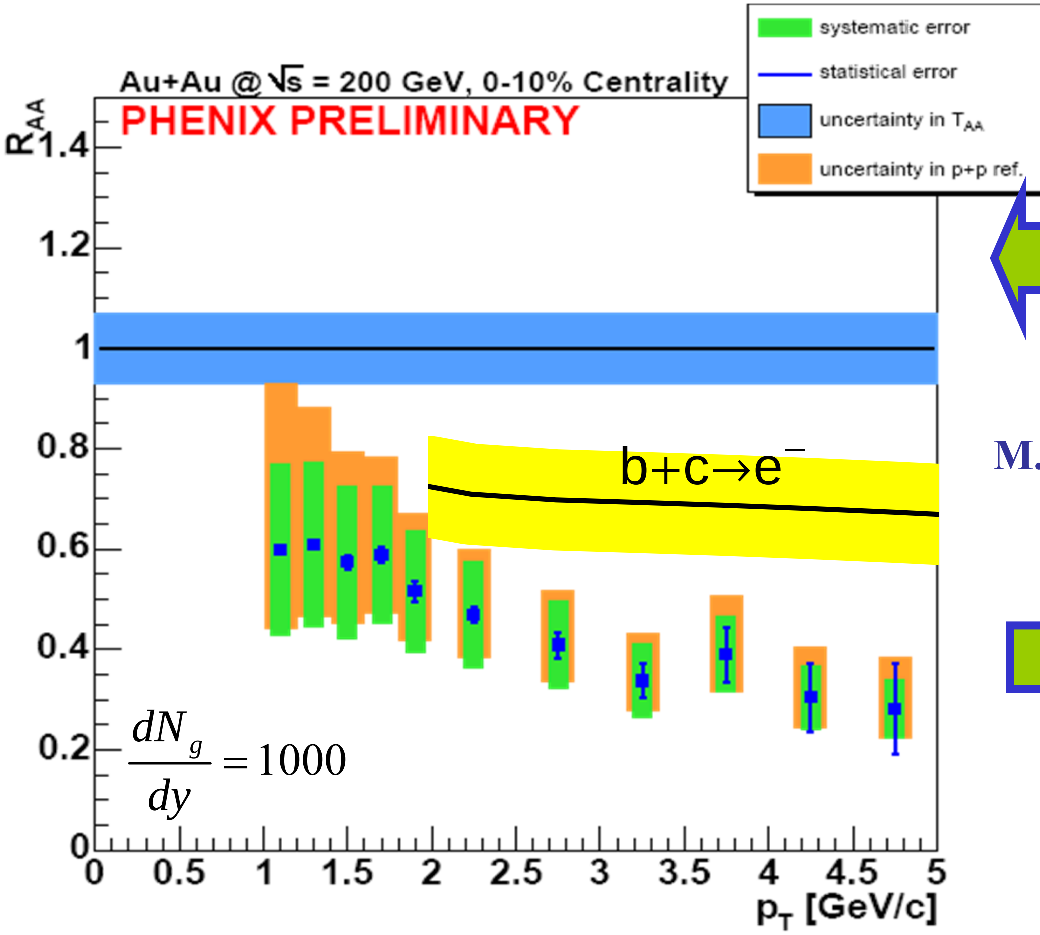


Considered to be negligible compared to radiative!

Single electron puzzle at RHIC



M. D. et al., Phys. Lett. B 632, 81 (2006)

Radiative energy loss predictions with $dN_g/dy=1000$

M. D. and M. Gyulassy, PRC 2003, PLB 2003, NPA 2004; M. D. PRC 2006;

Disagreement!

Radiative energy loss is not able to explain the single electron data as long as realistic parameter values are taken into account!

Does the radiative energy loss control the energy loss in QGP?

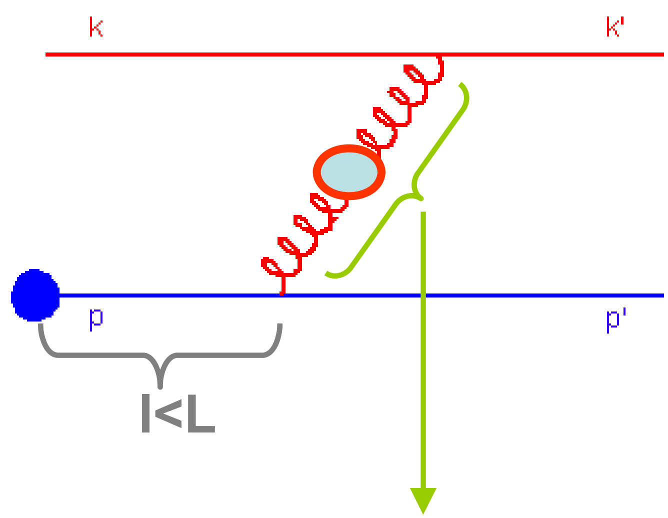
⇓

Is collisional energy loss also important?

Collisional energy loss in a finite size QCD medium

Consider a medium of size L in thermal equilibrium at temperature T.

The main order collisional energy loss is determined from:

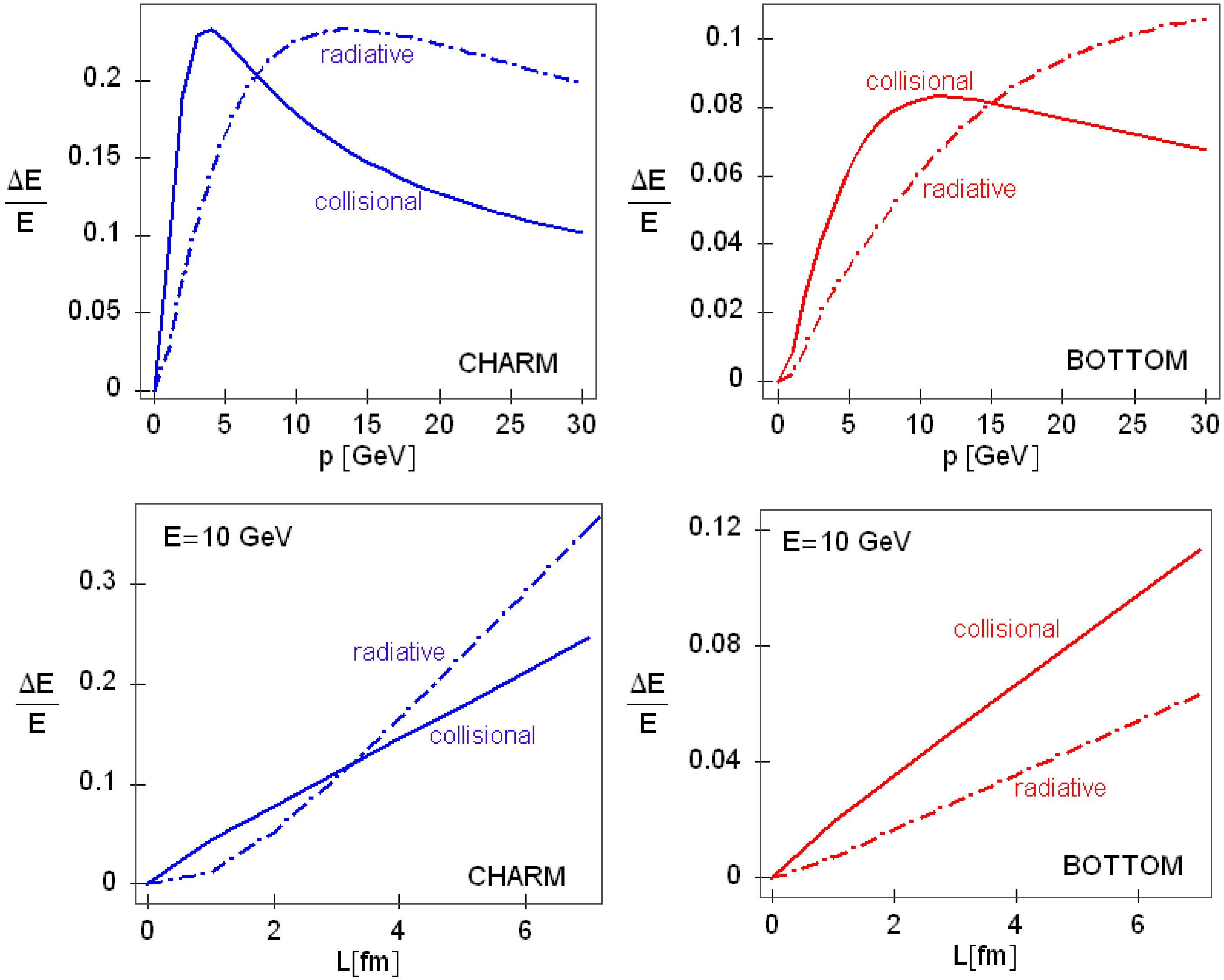


The effective gluon propagator:

$$D^{\mu\nu}(\omega, \vec{q}) = -P^{\mu\nu}\Delta_T(\omega, \vec{q}) - Q^{\mu\nu}\Delta_L(\omega, \vec{q})$$

M. D., Phys.Rev.C74:064907,2006

Collisional v.s. medium induced radiative energy loss



Collisional and radiative energy losses are comparable!

Non-zero collisional energy loss - a fundamental problem

Static QCD medium approximation (modeled by Yukawa potential).

→

With such approximation, collisional energy loss has to be exactly equal to zero!

↓

However, collisional and radiative energy losses are shown to be comparable.

←

Introducing collisional energy loss is necessary, but inconsistent with static approximation!

↓

Static medium approximation should not be used in radiative energy loss calculations!

→

Dynamical QCD medium effects have to be included!

Our goal

We want to compute both radiative and collisional energy loss in dynamical medium of thermally distributed massless quarks and gluons.

Why?

- **To address the applicability of static approximation in radiative energy loss computations.**
- **To compute collisional and radiative energy losses within a consistent theoretical framework.**

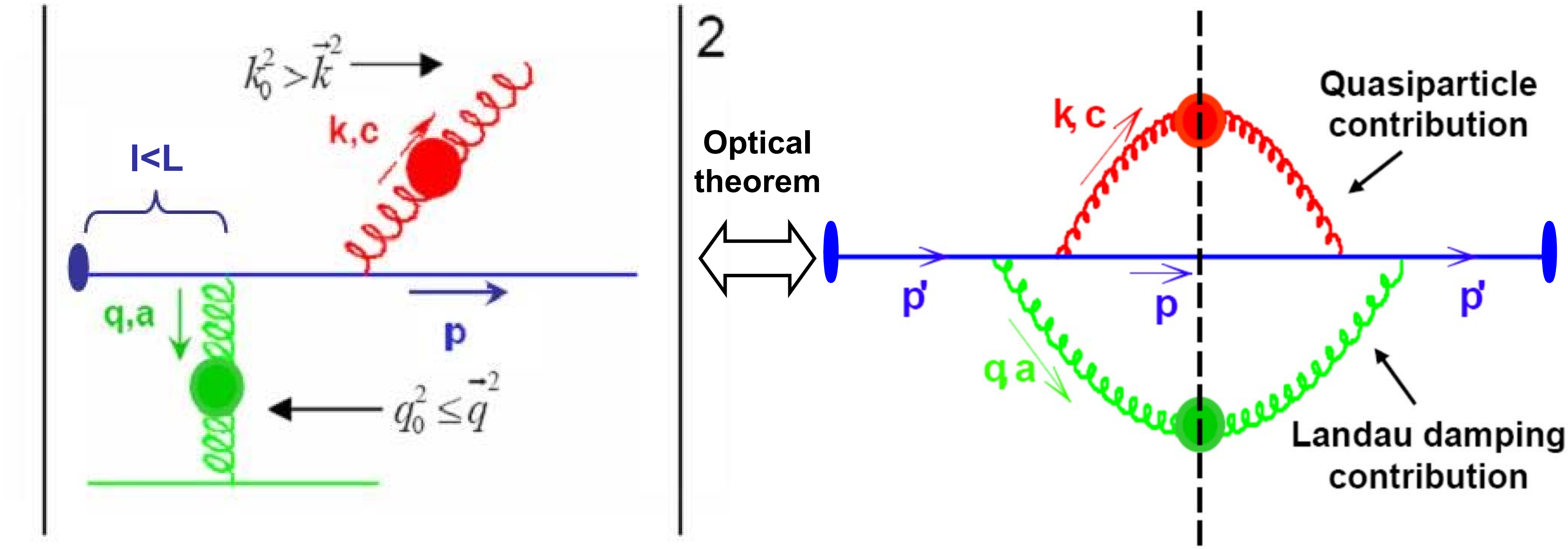
M. D., Phys.Rev.C80:064909,2009 (highlighted in APS physics).

M. D. and U. Heinz, Phys.Rev.Lett.101:022302,2008.

Radiative energy loss in a dynamical medium

We compute the medium induced radiative energy loss for a heavy quark to first (lowest) order in number of scattering centers.

To compute this process, we consider the radiation of one gluon induced by one collisional interaction with the medium.



We consider a medium of finite size L, and assume that the collisional interaction has to occur inside the medium.

The calculations were performed by using two Hard-Thermal Loop approach.

1-HTL gluon propagator:

$$iD^{\mu\nu}(l) = \frac{P^{\mu\nu}(l)}{l^2 - \Pi_T(l)} + \frac{Q^{\mu\nu}(l)}{l^2 - \Pi_L(l)}$$

Cut 1-HTL gluon propagator:

$$D^{>}_{\mu\nu}(l) = -(1+f(l_0))\Big(P_{\mu\nu}(l)\rho_T(l) + Q_{\mu\nu}(l)\rho_L(l)\Big) ,$$

$$\rho_{L,T}(l) = \underbrace{2\pi\,\delta(l^2 - \Pi_{T,L}(l))}_{\text{Radiated gluon}} - \underbrace{2\,\mathrm{Im}\left(\frac{1}{l^2 - \Pi_{T,L}(l)}\right)\theta(1 - \frac{l_0^2}{\vec{\mathrm{l}}^2})}_{\text{Exchanged gluon}}$$

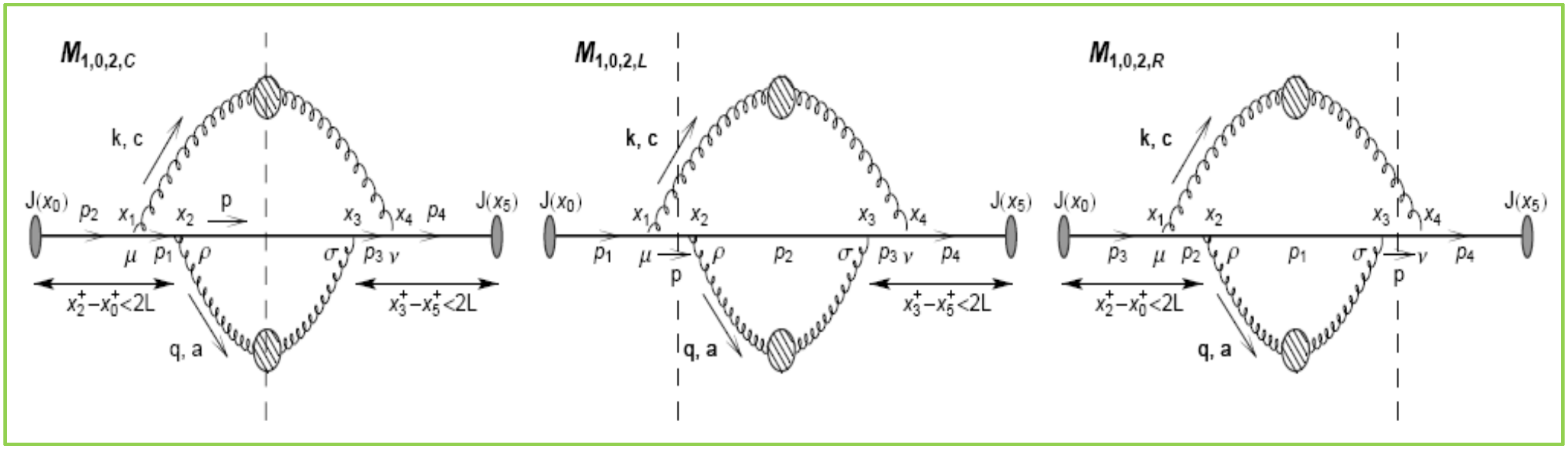
For radiated gluon, cut 1-HTL gluon propagator can be simplified to (M.D. and M. Gyulassy, PRC 68, 034914 (2003).

$$D^{>}_{\mu\nu}(k) \approx -2\pi \frac{P_{\mu\nu}(k)}{2\omega}\delta(k_0 - \omega) \qquad \omega \approx \sqrt{\vec{\mathrm{k}}^2 + m_g^2}\,;\; m_g \approx \mu/\sqrt{2}$$

For exchanged gluon, cut 1-HTL gluon propagator cannot be simplified, since both transverse (magnetic) and longitudinal (electric) contributions will prove to be important.

$$D^{>}_{\mu\nu}(q) = \theta(1 - \frac{q_0^2}{\vec{\mathrm{q}}^2})\,(1 + f(q_0))\,2\,\mathrm{Im}\left(\frac{P_{\mu\nu}(q)}{q^2 - \Pi_T(q)} + \frac{Q_{\mu\nu}(q)}{q^2 - \Pi_L(q)}\right)$$

More than one cut of a Feynman diagram can contribute to the energy loss in finite size dynamical QCD medium:



These terms interfere with each other, leading to the nonlinear dependence of the jet energy loss.

M. D., Phys.Rev.C80:064909,2009 (highlighted in APS physics).

We calculated all the relevant diagrams that contribute to this energy loss

Each individual diagram is infrared divergent, due to the absence of magnetic screening!

The divergence is naturally regulated when all the diagrams are taken into account.

So, all 24 diagrams have to be included to obtain sensible result.

$$\frac{\Delta E_{\rm dyn}}{E} = \frac{C_R\alpha_s}{\pi}\,\frac{L}{\lambda_{\rm dyn}}\int dx\,\frac{d^2\mathbf{k}}{\pi}\,\frac{d^2\mathbf{q}}{\pi}\,\frac{\mu^2}{\mathbf{q}^2(\mathbf{q}^2+\mu^2)}\left(1-\frac{\sin\frac{(\mathbf{k}+\mathbf{q})^2+\chi}{xE^+}L}{\frac{(\mathbf{k}+\mathbf{q})^2+\chi}{xE^+}L}\right)$$
$$\times\, 2\,\frac{(\mathbf{k}+\mathbf{q})}{(\mathbf{k}+\mathbf{q})^2+\chi}\left(\frac{(\mathbf{k}+\mathbf{q})}{(\mathbf{k}+\mathbf{q})^2+\chi}-\frac{\mathbf{k}}{\mathbf{k}^2+\chi}\right),$$

M. D., Phys.Rev.C80:064909,2009 (highlighted in APS physics).

Finite magnetic mass

The dynamical energy loss formalism is based on HTL perturbative QCD, which requires zero magnetic mass.

↓

However, different non-perturbative approaches show a **non-zero magnetic mass at RHIC and LHC**.

↓

Can magnetic mass be consistently included in the dynamical energy loss calculations?

Generalization of radiative jet energy loss to finite magnetic mass

$$\frac{\Delta E_{\rm dyn}}{E} = \frac{C_R\alpha_s}{\pi}\frac{L}{\lambda_{\rm dyn}}\int dx \frac{d^2k}{\pi}\frac{d^2q}{\pi}\frac{\mu^2}{q^2(q^2+\mu^2)} \times 2\frac{(k+q)}{(k+q)^2+\chi}\left(\frac{(k+q)}{(k+q)^2+\chi}-\frac{k}{k^2+\chi}\right)\left(1-\frac{\sin\frac{(k+q)^2+\chi}{xE^+}L}{\frac{(k+q)^2+\chi}{xE^+}L}\right)$$

zero magnetic mass

From our analysis, only this part gets modified.

Finite magnetic mass: $\frac{\mu_E^2-\mu_M^2}{(q^2+\mu_E^2)(q^2+\mu_M^2)}$, **where** $0.4 \leq \frac{\mu_M}{\mu_E} \leq 0.6$.

M.D. and M. Djordjevic, Phys.Lett.B709:229,2012

Dynamical energy loss - summary

Computed both collisional and radiative energy loss, in a finite size QCD medium, composed of dynamical scatterers.

M. D. PRC 80:064909 (2009), M. D. and U. Heinz, PRL 101:022302 (2008).

Finite magnetic mass effects

M. D. and M. Djordjevic, PLB 709:229 (2012)

Includes running coupling

M. D. and M. Djordjevic, PLB 734, 286 (2014).

State of the art energy loss formalism in a dynamical finite size QCD medium.

Numerical importance of different effects addressed in

B. Blagojevic and M.D, J.Phys. G42 (2015) 7, 075105 (highlighted in LabTalk)

Numerical procedure

- **Light flavor production** Z.B. Kang, I. Vitev, H. Xing, PLB 718:482 (2012)
- **Heavy flavor production** M. Cacciari et al., JHEP 1210, 137 (2012)
- **Path-length fluctuations** A. Dainese, EPJ C33:495,2004.
- **Multi-gluon fluctuations**

  M. Gyulassy, P. Levai, I. Vitev, PLB 538:282 (2002).
- **DSS and KKP fragmenation for light flavor**

  D. de Florian, R. Sassot, M. Stratmann, PRD 75:114010 (2007)

  B. A. Kniehl, G. Kramer, B. Potter, NPB 582:514 (2000)
- **BCFY and KLP fragmenation for heavy flavor**

  M. Cacciari, P. Nason, JHEP 0309: 006 (2003)
- **Decays of heavy mesons to single electron and J/ψ according to**

  M. Cacciari et al., JHEP 1210, 137 (2012)
- **Temperature T=304 MeV for LHC and T=221 MeV for RHIC.**

  M. Wilde, Nucl. Phys. A 904-905, 573c (2013) (ALICE Collab.)

  A. Adare *et al*., Phys. Rev. Lett. 104, 132301 (2010) (PHENIX Collab.)

Comparison with the experimental data

- **Provide joint predictions across diverse probes**

  charged hadrons, pions, kaons, D mesons,

  non-photonic single electorns, non-prompt J/ψ

  M. D. and M. Djordjevic, PLB 734, 286 (2014)

- **Concentrate on all centrality regions**

  M. D., M. Djordjevic and B. Blagojevic, PLB 737 298 (2014)

- **Provide predictions for the upcoming data**

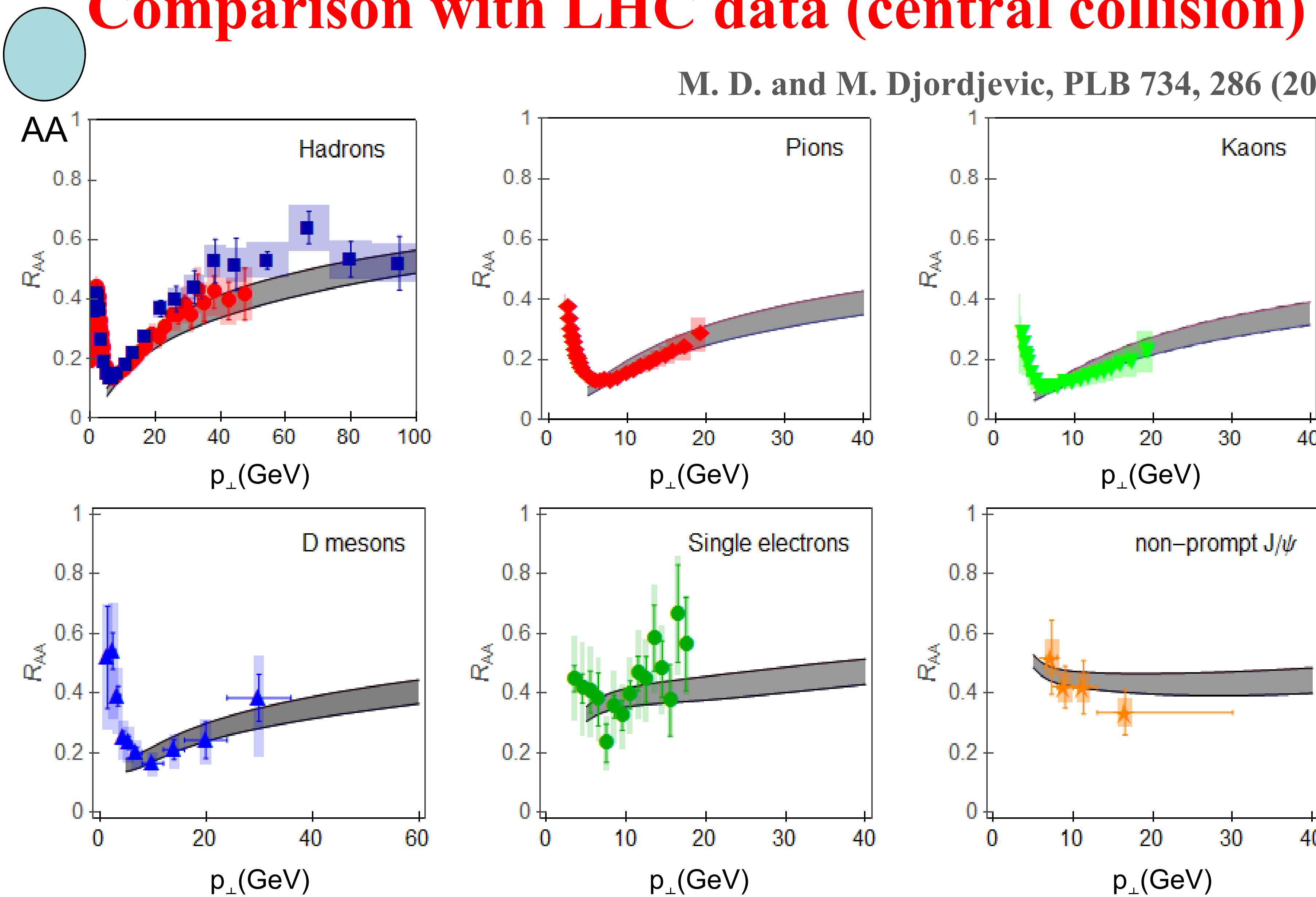
  M. D. and M. Djordjevic, Phys. Rev. C 92 (2015) 2, 024918

  M. D., B. Blagojevic and L. Zivkovic, Phys. Rev. C 94 (2016) 044908

- **All predictions generated**
  - **By the same formalism**
  - **With the same numerical procedure**
  - **No free parameters in model testing**

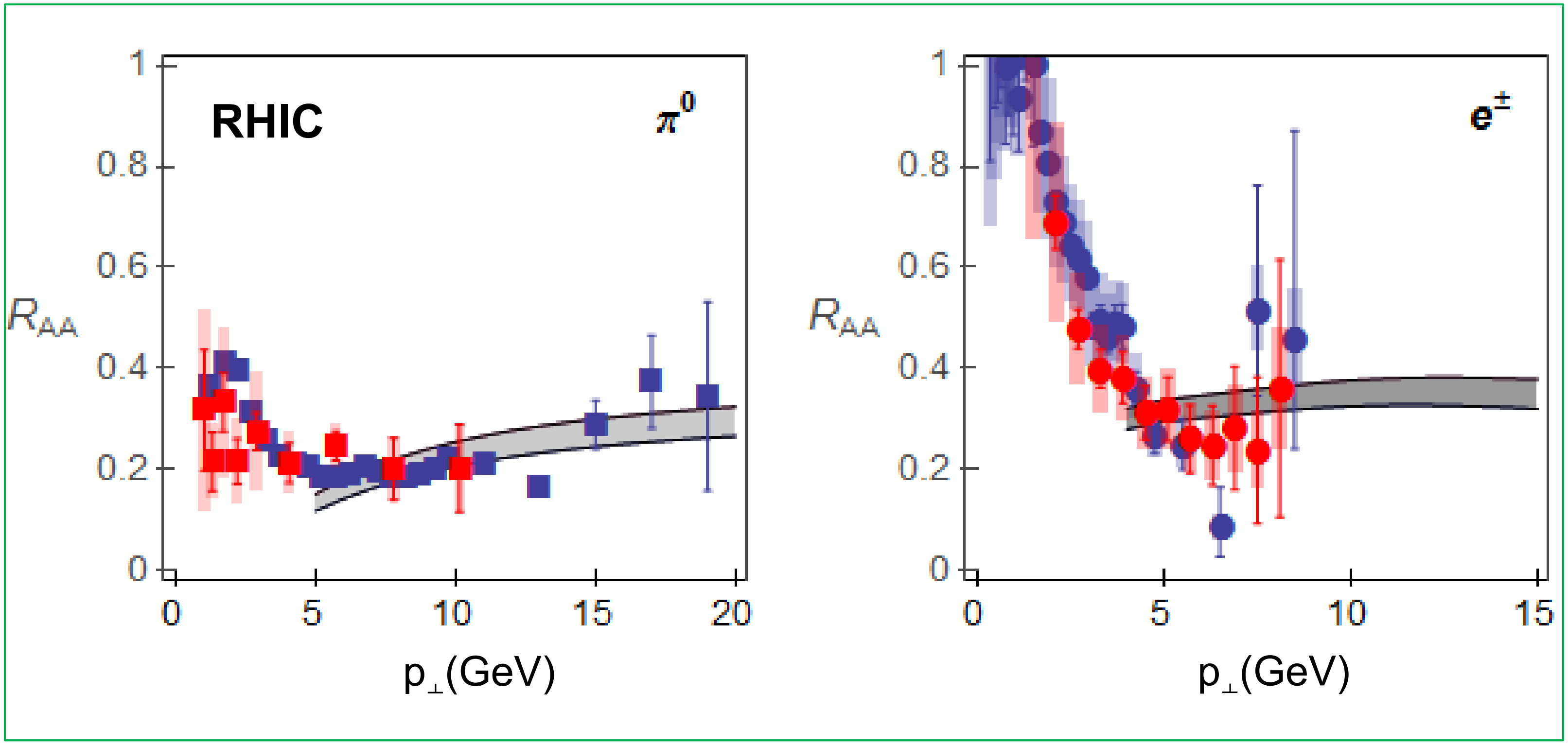
Comparison with LHC data (central collision)

M. D. and M. Djordjevic, PLB 734, 286 (2014)



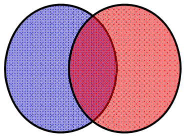
Very good agreement with diverse probes!

Comparison with RHIC data (central collisions)

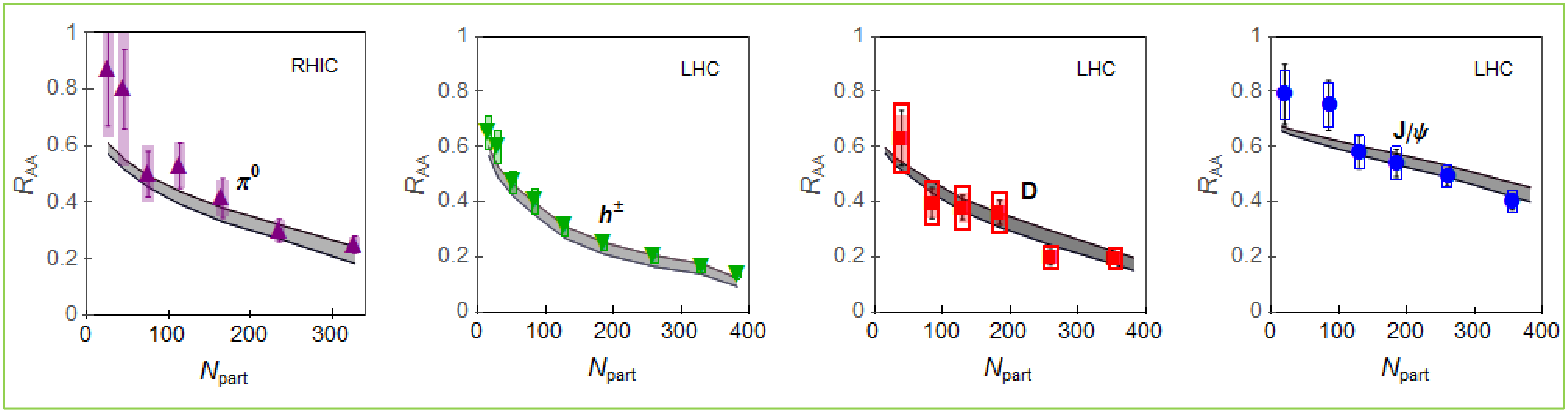


Very good agreement!

M.D. and M. Djordjevic, PRC 90, 034910 (2014)



R_{AA} *vs.* N_{part} for RHIC and LHC

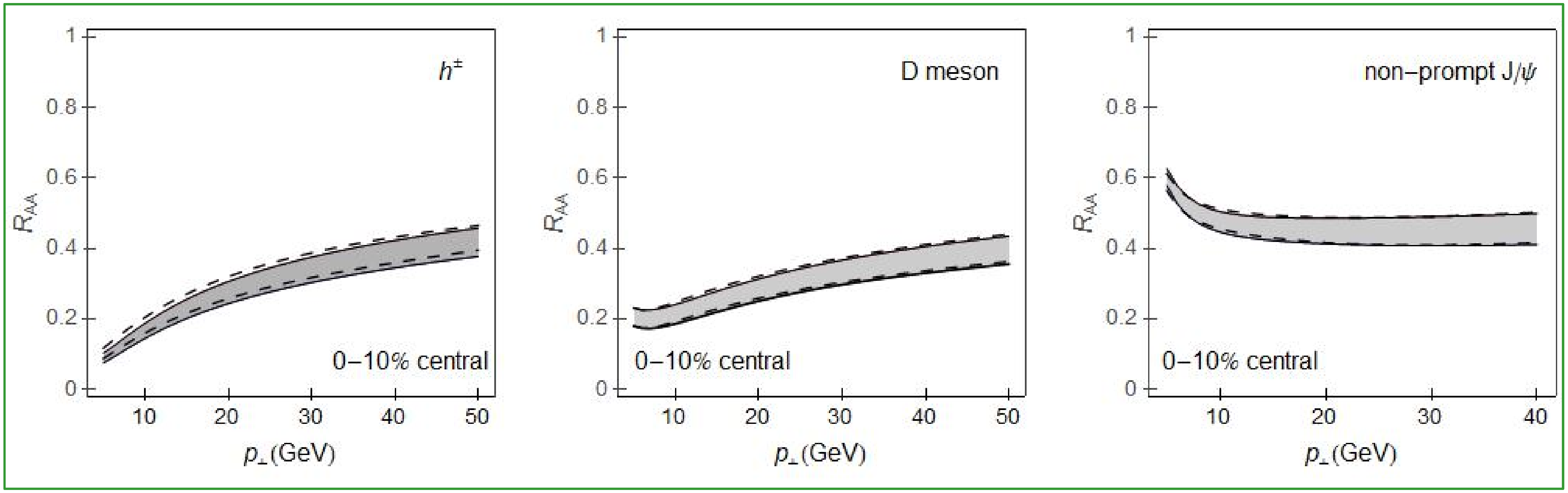


Excellent agreement for both RHIC and LHC and for the whole set of probes!

M. D., M. Djordjevic and B. Blagojevic, PLB 737 298 (2014)

5.02 TeV Pb+Pb at LHC

M. D. and M. Djordjevic, Phys. Rev. C 92 (2015) 2, 024918



The same suppression as at 2.76 TeV for all types of probes!

Confirmed by experimental data in July 2016!

Summary

Dynamical energy loss formalism.

Tested on angular averaged R_{AA} data

Largely not sensitive to the medium evolution.

Good agreement for wide range of probes, centralities and beam energies.

Can explain puzzling data.

Clear predictions for future experiments.

The dynamical energy loss formalism can well explain the jet-medium interactions in QGP.